

So, one has,

$$\frac{\partial \Psi}{\partial x} = \frac{i p_x}{\hbar} \underbrace{A e^{i(p_x x - E t)/\hbar}}_{\Psi(x,t)}$$

$$\Rightarrow \boxed{\hat{p}_x = -i\hbar \frac{\partial}{\partial x}}, \text{ Similarly } \boxed{E = i\hbar \frac{\partial}{\partial t}}$$

$$\text{So, } \hat{T} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \left\{ \hat{T} = \frac{\hat{p}_x^2}{2m} \right\}$$

Eigenvalue Equation: معادلة القيمة الذاتية

For some function Ψ :

$$\hat{A} \Psi = \alpha \Psi$$

$$\text{or } \boxed{\hat{L} \Psi = \lambda \Psi}$$

where α is a complex number. In this case, if Ψ is a member of class of physically meaningful functions, it is an eigenfunction of the operator \hat{A} . The number (α) is called the eigenvalue of \hat{A} associated with the eigenfunction or state function Ψ .

Example 3 Show that the wave function $\psi = Nxe^{-\frac{x^2}{2}}$ is an eigen function of $H = -\frac{d^2}{dx^2} + x^2$

Also, find out the normalization constant N and the eigen value.

$$\begin{aligned}
 H\psi &= \left(-\frac{d^2}{dx^2} + x^2\right) Nxe^{-\frac{x^2}{2}} \\
 &= Nxe^{-\frac{x^2}{2}} - N \frac{d}{dx} \left[e^{-\frac{x^2}{2}} - xe^{-\frac{x^2}{2}} \right] \\
 &= 3Nxe^{-\frac{x^2}{2}} \\
 &= 3\psi
 \end{aligned}$$

Hence the eigenvalue is 3.

From normalization condition,

$$\begin{aligned}
 N^2 \int_{-\infty}^{+\infty} x^2 e^{-x^2} dx &= 1 \\
 2N^2 \int_0^{\infty} x^2 e^{-x^2} dx &= 1
 \end{aligned}$$

$$\begin{aligned}
 x^2 = t &\Rightarrow dx = \frac{1}{2} t^{-\frac{1}{2}} dt \\
 x = t^{\frac{1}{2}} &\Rightarrow dx = \frac{1}{2} t^{-\frac{1}{2}} dt
 \end{aligned}$$

$$\begin{aligned}
 \Rightarrow 2N^2 \int_0^{\infty} t e^{-t} \cdot \frac{1}{2} t^{\frac{1}{2}} dt &= 1 \\
 2N^2 \cdot \frac{1}{2} \int_0^{\infty} t^{\frac{1}{2}} e^{-t} dt &= 1
 \end{aligned}$$

-31-

$$\Gamma(p) = \int_0^{\infty} x^{p-1} e^{-x} dx \quad \text{Gamma function}$$

$$\Gamma\left(\frac{3}{2}\right) = \int_0^{\infty} x^{\frac{1}{2}} e^{-x} dx ; \quad \Gamma(p+1) = p\Gamma(p)$$

$$\Gamma\left(\frac{3}{2}\right) = \Gamma\left(\frac{1}{2}+1\right) = \frac{1}{2}\Gamma\left(\frac{1}{2}\right)$$

$$= \frac{\sqrt{\pi}}{2}$$

$$\Rightarrow N^2 \cdot \frac{\sqrt{\pi}}{2} = 1$$

$$N^2 = \frac{2}{\sqrt{\pi}} \Rightarrow N = \left(\frac{2}{\sqrt{\pi}}\right)^{1/2}$$

Hermitian Operators and their properties

الـؤثرات الـهرميتية

An operator \hat{A} or \hat{F} , for example, associated with a dynamical variable is said to be Hermitian if its average or expectation value in any eigenstate ψ is real.

The average value of an operator \hat{A} for the normalized eigenfunction ψ is expressed as

$$\langle A \rangle = \int \psi^* A \psi dx$$

Or the Condition for an operator \hat{A} to be Hermitian is, تعريف المؤثر الهرميتي

$$\int_{\text{all space}} \psi_1^* \hat{A} \psi_2 dv = \int \psi_2 (\hat{A} \psi_1)^* dv$$

Example / Show that the momentum operator $\hat{p}_x = -i\hbar \frac{\partial}{\partial x}$ is a Hermitian operator.

$$\int_{-\infty}^{+\infty} \psi_1^* \hat{A} \psi_2 dv = \int_{-\infty}^{+\infty} \psi_2 (\hat{A} \psi_1)^* dv$$

$$\int_{-\infty}^{+\infty} \psi_1^* (-i\hbar \frac{\partial}{\partial x}) \psi_2 dx = -i\hbar \int_{-\infty}^{+\infty} \psi_1^* \frac{\partial \psi_2}{\partial x} dx$$

$$\int u dv = uv - \int v du$$

let $u = \psi_1^*$ $du = \frac{\partial \psi_1^*}{\partial x} dx$

$$\frac{\partial \psi_2}{\partial x} dx = dv \Rightarrow v = \psi_2$$

$$\Rightarrow -i\hbar \int_{-\infty}^{+\infty} \psi_1^* \frac{\partial \psi_2}{\partial x} dx = -i\hbar \left[\psi_1^* \psi_2 \right]_{-\infty}^{+\infty} - \int_{-\infty}^{+\infty} \psi_2 \frac{\partial \psi_1^*}{\partial x} dx$$

$$= +i\hbar \int_{-\infty}^{+\infty} \psi_2 \frac{\partial \psi_1^*}{\partial x} dx$$

$$= \int_{-\infty}^{+\infty} \psi_2 (i\hbar \frac{\partial}{\partial x}) \psi_1^* dx$$