

properties of Hermitian operators :

1. They have real eigenvalues.
2. Any set of mutually orthogonal eigenfunctions is linearly independent.

proof of property [1] : اثبات اولی

Let (f) be an eigenvalue of the Hermitian operator \hat{A} in the state described by the normalized wavefunction ψ . The eigenvalue equation is :

$$\hat{A}\psi = f\psi \quad \text{--- (1)}$$

its complex conjugate is :

$$\hat{A}^* \psi^* = f^* \psi^* \quad \text{--- (2)}$$

multiplying eq. by ψ^* , and eq. (2) by ψ , one has :

$$\psi^* \hat{A} \psi = \psi^* f \psi \quad \text{--- (3)}$$

$$\psi \hat{A}^* \psi^* = \psi f^* \psi^* \quad \text{--- (4)}$$

by integrating the eqs. (3), (4), respectively, one obtains:

$$\int \psi^* A \psi dv = f \int \psi^* \psi dv$$

also,

$$\int \psi A^* \psi^* dv = f^* \int \psi \psi^* dv$$

where $\psi_1 = \psi_2 = \psi$

But $\int \psi^* \psi dv = 1$ for a normalized wavefunction ψ ,

$$f - f^* = 0 \Rightarrow \boxed{f^* = f}$$

which is possible only if (f) is a real number. Therefore, eigenvalue of Hermitian operator is real.

proof of property (2):

Assume ψ_1, ψ_2 are two ~~eigenfunctions~~ eigenfunctions for the Hermitian operator A and their eigenvalues f_1, f_2 respectively,

We apply the eigenvalue equation, that is

$$\left. \begin{aligned} \hat{A}\psi_1 &= E_1 \psi_1 && \text{--- (1)} \\ \hat{A}\psi_2 &= E_2 \psi_2 && \text{--- (2)} \end{aligned} \right\}$$

the complex conjugate of eq (1) :

$$\hat{A}^* \psi_1^* = E_1^* \psi_1^*$$

from the property [1], the last equation becomes

$$\hat{A} \psi_1^* = E_1 \psi_1^* \text{ --- (3)}$$

Multiplying eq (3) by ψ_2 and eq (2) by ψ_1^* that is

$$\left. \begin{aligned} \psi_2 \hat{A} \psi_1^* &= \psi_2 E_1 \psi_1^* && \text{--- (4)} \\ \psi_1^* \hat{A} \psi_2 &= \psi_1^* E_2 \psi_2 && \text{--- (5)} \end{aligned} \right\}$$

by integrating the eqs. (4), and (5) and subtracting one has

$$\int \psi_2 \hat{A} \psi_1^* dv - \int \psi_1^* \hat{A} \psi_2 dv = (E_1 - E_2) \int \psi_2 \psi_1^* dv$$

all space

According to definition of Hermitian operator

$$\Rightarrow (E_1 - E_2) \int \psi_2 \psi_1^* dv = 0$$

$$\left. \begin{aligned} E_1 \neq E_2 &\Rightarrow \int \psi_2 \psi_1^* dv = 0 \end{aligned} \right\}$$

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i.e., ψ_1, ψ_2 are linearly independent
 (Orthogonal wavefunctions); E_1, E_2 are
 different eigenvalues ($E_1 \neq E_2$)

Ex

Example: Show that the operator
 $-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$ is a Hermitian
 operator. مؤثر الطاقة الحركية

Sol. التعرف

$$\int_{-\infty}^{+\infty} \psi_1^* A \psi_2 dx \stackrel{?}{=} \int_{-\infty}^{+\infty} \psi_2 (A \psi_1)^* dx$$

the left-hand side:

$$-\frac{\hbar^2}{2m} \int_{-\infty}^{+\infty} \psi_1^* \frac{\partial^2}{\partial x^2} \psi_2 dx = ?$$

by integration by parts, one obtains:

$$-\frac{\hbar^2}{2m} \int_{-\infty}^{+\infty} \psi_1^* \frac{\partial^2}{\partial x^2} \psi_2 dx$$

let $\psi_1^* = u \Rightarrow du = \frac{\partial \psi_1^*}{\partial x} dx$; $dv = \frac{\partial^2}{\partial x^2} \psi_2 dx$

$$v = \frac{\partial}{\partial x} \psi_2$$

$$\Rightarrow \int u dv = uv - \int v du$$

$$\frac{-\hbar^2}{2m} \int \psi^* \frac{\partial^2}{\partial x^2} \psi dx = \frac{\hbar^2}{2m} \left[\psi^* \frac{\partial \psi}{\partial x} \right]_{-\infty}^{\infty} + \frac{\hbar^2}{2m} \int \frac{\partial \psi^*}{\partial x} \frac{\partial \psi}{\partial x} dx$$

$$= \frac{\hbar^2}{2m} \int \frac{\partial \psi^*}{\partial x} \frac{\partial \psi}{\partial x} dx$$

Now,

$$u = \frac{\partial \psi^*}{\partial x}$$

$$du = \frac{\partial^2 \psi^*}{\partial x^2} dx$$

$$dv = \frac{\partial}{\partial x} \psi dx \Rightarrow dv = \psi$$

$$\int u dv = uv - \int v du$$

$$\frac{\hbar^2}{2m} \int \frac{\partial \psi^*}{\partial x} \psi \frac{\partial \psi^*}{\partial x} dx = \left[\frac{\partial \psi^*}{\partial x} \psi \right]_{-\infty}^{\infty} - \frac{\hbar^2}{2m} \int \psi \frac{\partial^2 \psi^*}{\partial x^2} dx$$

$$\int_{-\infty}^{\infty} \frac{-\hbar^2}{2m} \psi^* \frac{\partial^2}{\partial x^2} \psi dx = -\frac{\hbar^2}{2m} \int \psi \left(\frac{\partial^2}{\partial x^2} \psi^* \right) dx$$

اذا مؤثر الطاقة
فواز هيرميتي

Expectation value : القيمة المتوقعة

The average of dynamical quantity (observable) is given by

$$\langle A \rangle = \frac{\int \psi^* (A \psi) dx}{\int \psi^* \psi dx} \quad (1)$$

if ψ is a normalized wavefunction, then $\int \psi^* \psi dv = 1$ "normalization condition"

So, eq. (1) becomes :

$$\langle A \rangle = \int_{\text{all space}} \psi^* A \psi dv$$

Now, if the wavefunction ψ is ~~also~~ an eigen one ~~of~~ of the operator \hat{A} with eigenvalue f , then :

$$\langle A \rangle = f \int \psi^* \psi dv ; \quad \int \psi^* \psi dv = 1$$

$\therefore \langle A \rangle = f$ i.e., the expectation value equals the eigen value.

* To evaluate the expectation value of

P_x , one obtains :

$$\langle P_x \rangle = \int_{-\infty}^{+\infty} \psi^* P_x \psi dx$$

$$\text{[or]} \quad \langle P_x \rangle = -i\hbar \int \psi^* \frac{\partial}{\partial x} \psi dx$$

and so on, one can obtain the expectation value

for any observable.

H.W. دایب بی Show that the wavefunction

$$\psi(x) = b e^{-ax} ; a, b \text{ are constants}$$

is an eigen wavefunction of the operator $\frac{d^2}{dx^2}$.

Uncertainty Relations میشو، تویس

The uncertainty in the position Δx can be expressed as,

$$\Delta x = \left\{ \langle x^2 \rangle - \langle x \rangle^2 \right\}^{1/2}$$

Also, the uncertainty in the momentum coordinate Δp is expressed as,

$$\Delta p = \left\{ \langle p^2 \rangle - \langle p \rangle^2 \right\}^{1/2}$$

for example,

$$\langle p^2 \rangle = \int_{-\infty}^{\infty} \psi^* p^2 \psi dx$$

$$p^2 = -\hbar^2 \frac{\partial^2}{\partial x^2} \Rightarrow \langle p^2 \rangle = -\hbar^2 \int_{-\infty}^{\infty} \psi^* \frac{\partial^2}{\partial x^2} \psi dx$$